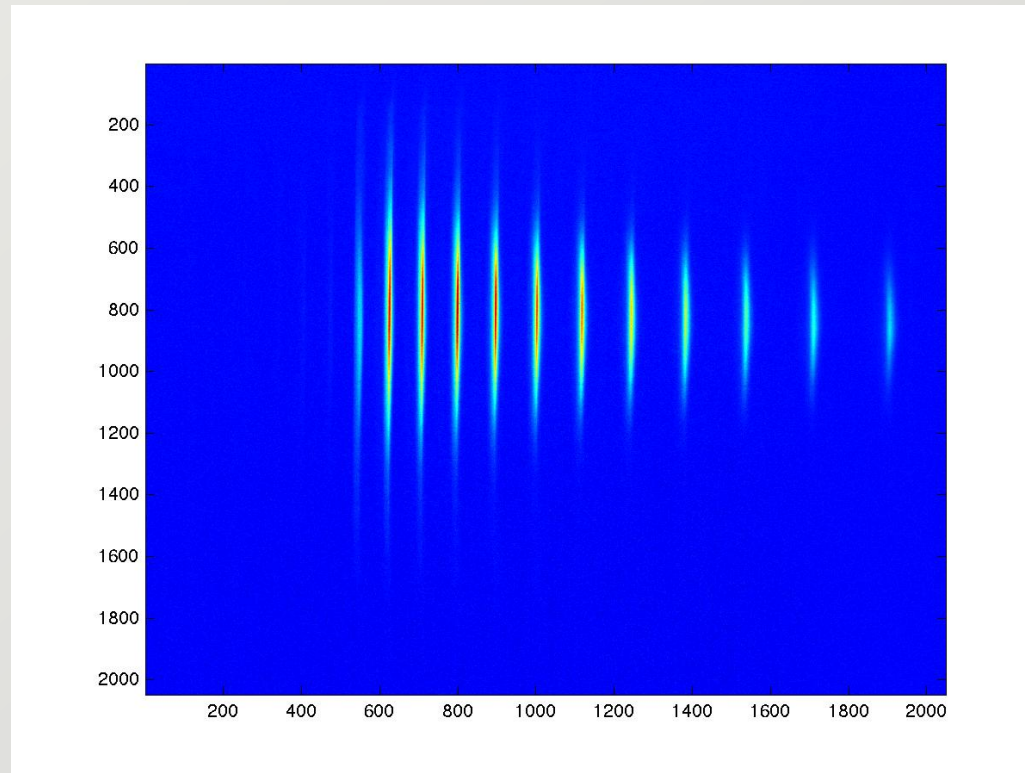


# Coherent extreme ultra violet light sources using highly efficient High Harmonic Generation



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# Alternative EUV sources



Synchrotron/FEL: large scale facility is needed



Plasma sources: limited by atomic line emission



Plasma sources: high divergence ( $4\pi$ )

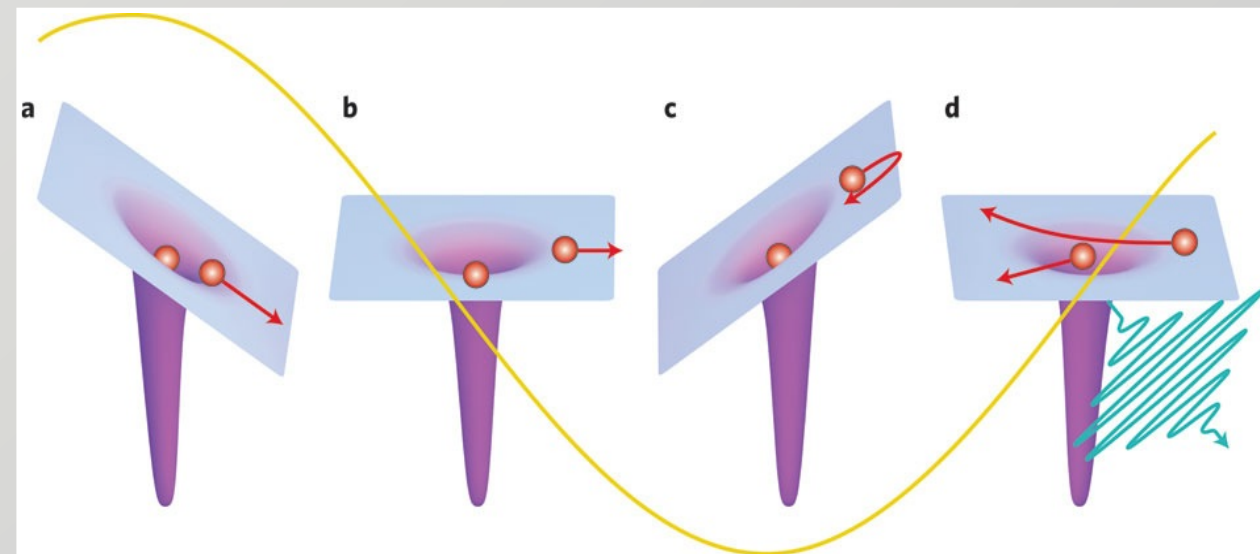


monochromatic sources



Could HHG be an alternative source?

Midorikawa, Nature Photonics 5, 640–641 (2011)





# Outline

1. Mechanism of High Harmonic Generation – HHG
2. Using different concepts of HHG
3. Field of application – XUV Coherence tomography



# Properties of HHG



Pulse duration in attosecond regime



Laser-like radiation



driven by small scale lab-based fs-lasers



intrinsically broadband



efficiency up to  $10^{-5}$





# Mechanism of High Harmonic Generation



Illuminate an atom with an intense laser field –  
what happens?



not a simple field ionization:  $E_{\text{photon}} > U_{\text{atom}}$



for NIR-laser and gases:  $E_{\text{photon}} \ll U_{\text{atom}}$



High non-linearity:  $N \cdot E_{\text{photon}} = U_{\text{atom}}$

Example	photon energy	Number of photons
Hydrogen ( $U=13.6$ eV)	800 nm (1.55 eV)	>8
Helium ( $U=24.58$ eV)	800 nm (1.55 eV)	>15



# Non-perturbative: Three-step model

Corkum, PRL 71 13, 1994-1997 (1993)



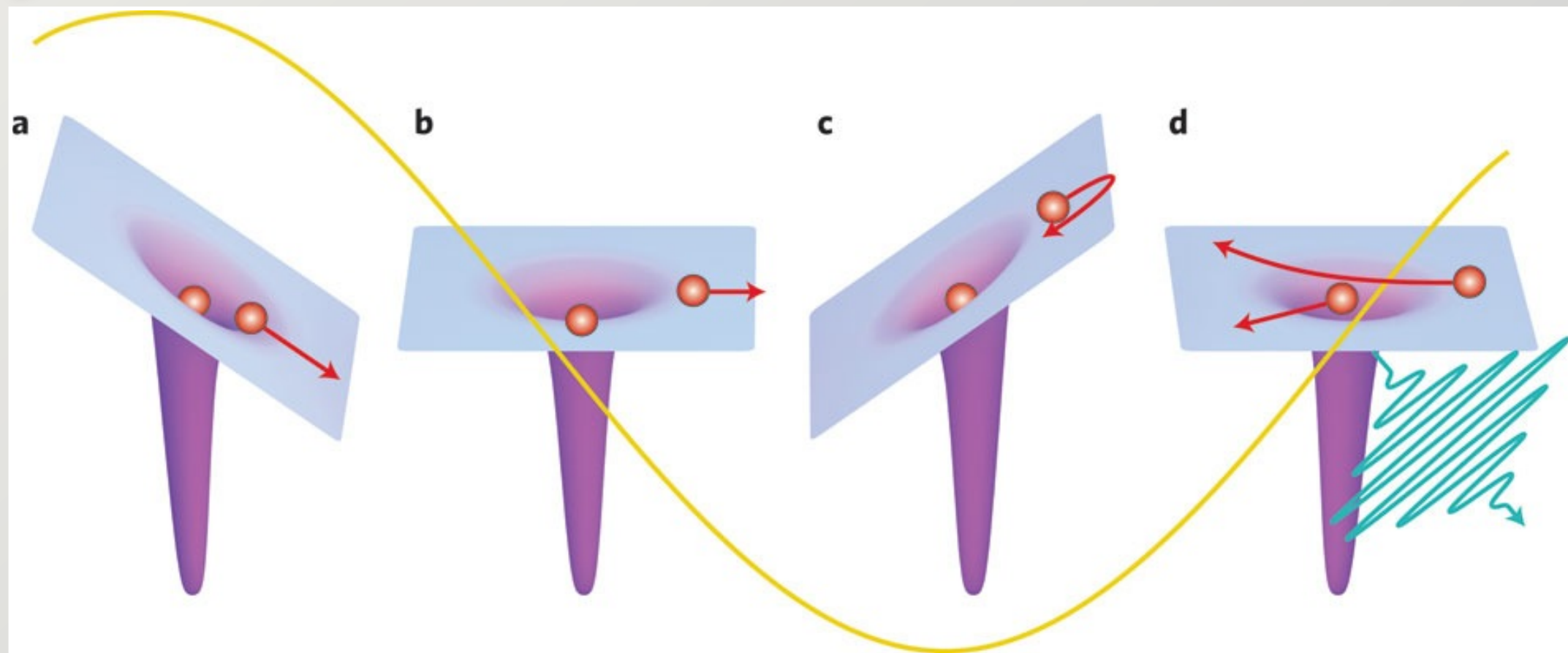
a. Ionization



b./c. Propagation in the laser field



d. Interaction with parent ion



Midorikawa, Nature Photonics 5, 640–641  
(2011)

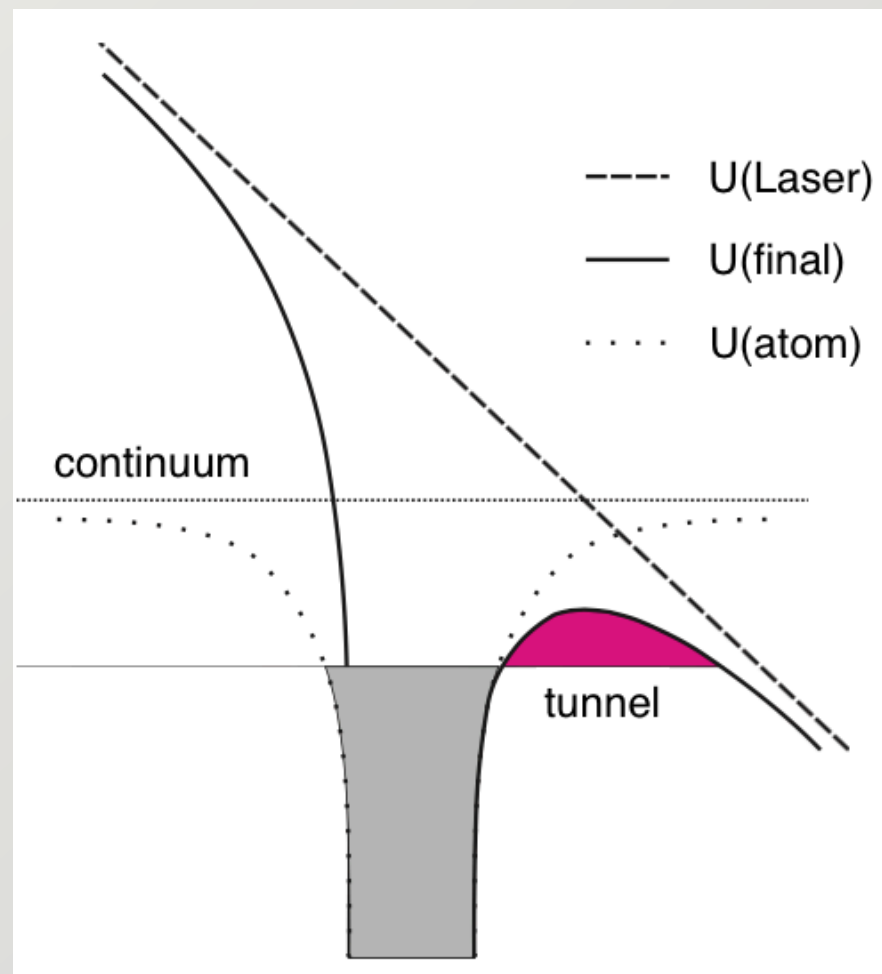


# Ionization



Potential deformation in linear fields

$$U_{\text{final}}(\vec{r}, t) = U_{\text{atom}} + U_{\text{laser}} = -\frac{1}{4\pi\epsilon_0} \frac{Ze^2}{|\vec{r}|} - e\vec{E}(\vec{r}, t)\vec{r}$$



Tunneling through the atomic potential



# Propagation in the laser field

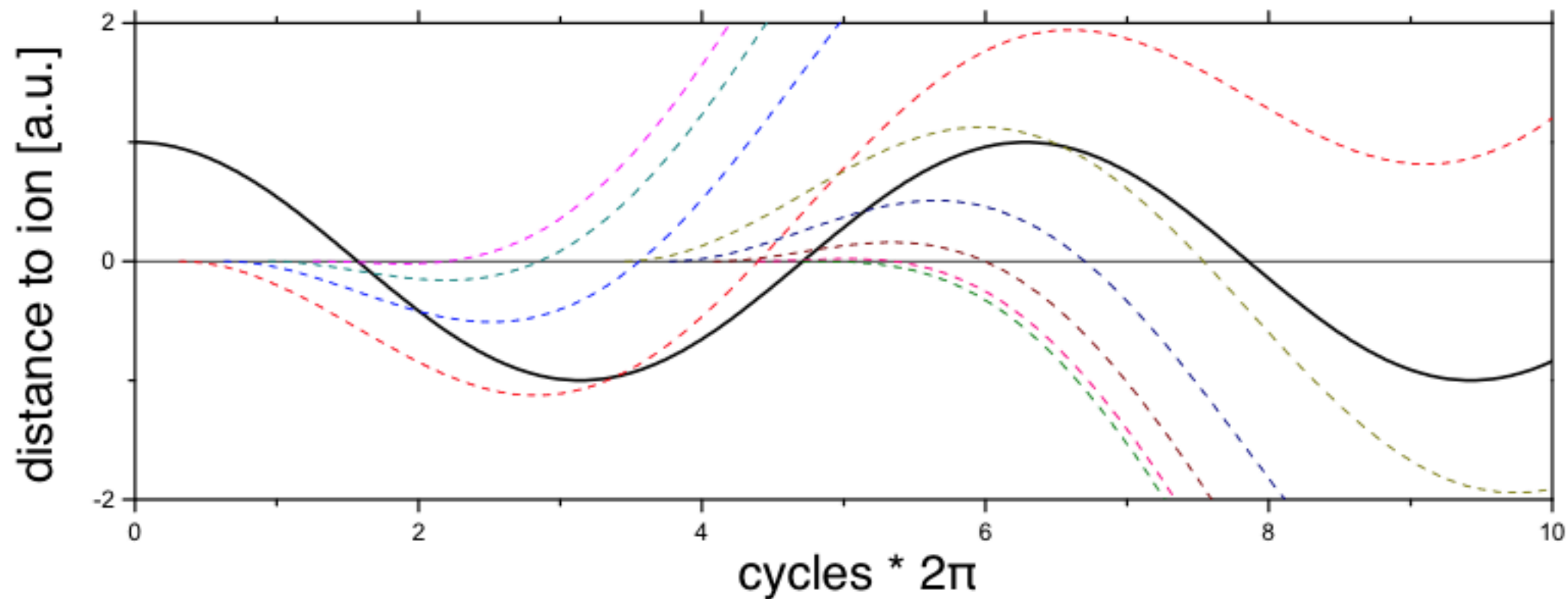


Electron follows the field



Canonical momentum with vector potential  $\vec{A}$

$$\vec{p}(t_0) + e\vec{A}(t_0) = \vec{p}(t_1) + e\vec{A}(t_1)$$



Can return to the parent ion





# Recombination



Return energy

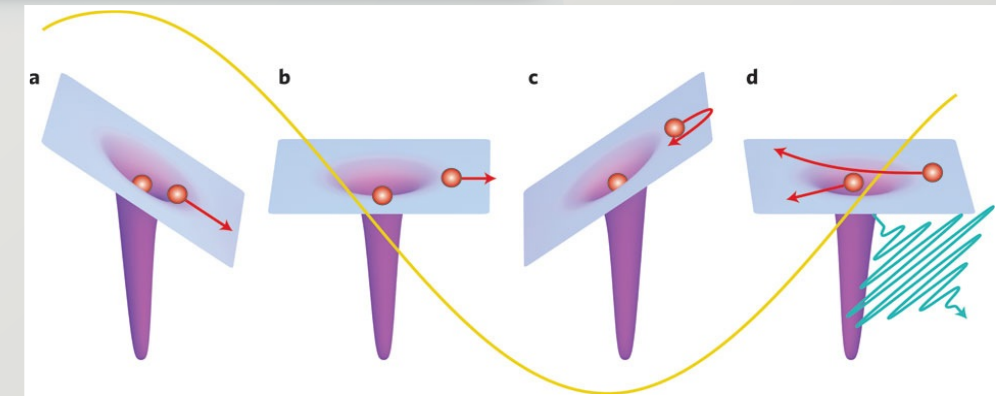
$$E_{\text{return}} = \frac{e^2}{2m} (A(t_1) - A(t_0))^2 \leq 3.17 U_P$$



ponderomotive potential = kinetic energy of the laser field



Cut-off energy:  $\hbar\omega_{\text{CO}} = I_P + 3.17 U_P$   
 $U_P \approx 10^{-13} I \lambda^2$



photon energy	intensity	Ponderomotive Potential	Cut-off
800 nm (1.55 eV)	$10^{14} \text{ W/cm}^2$	$U_P = 6.4 \text{ eV}$	$E = 36 \text{ eV (@Ar)}, \lambda = 34 \text{ nm}$
1800 nm (0.69 eV)	$10^{14} \text{ W/cm}^2$	$U_P = 32.4 \text{ eV}$	$E = 118 \text{ eV (@Ar)}, \lambda = 10.6 \text{ nm}$



# beam characteristics



periodic process every laser half-cycle



Frequency comb

$$E_{\text{harm}} = N \cdot E_0$$



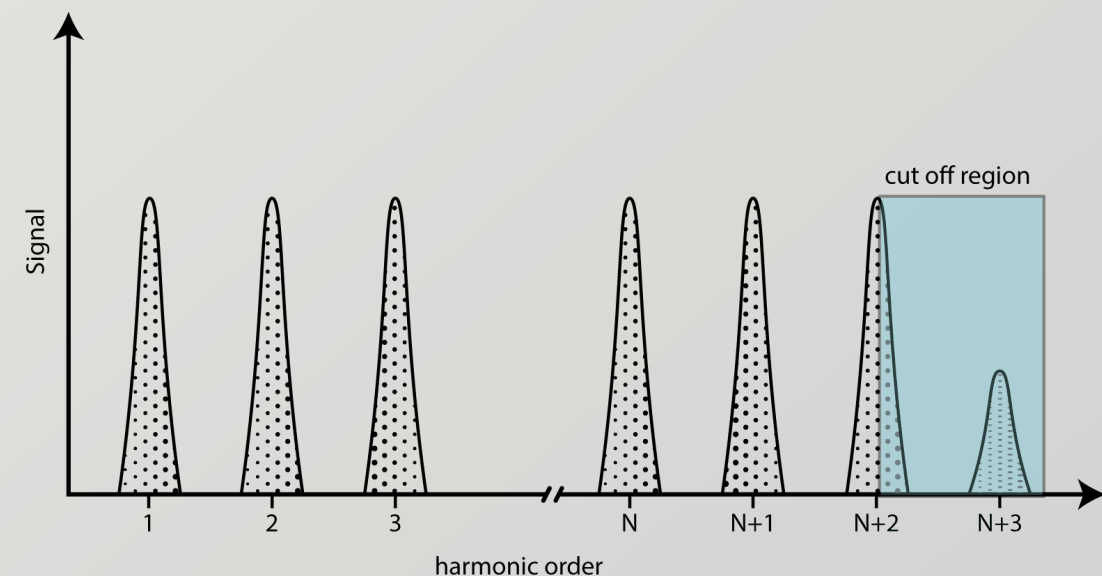
efficiency up to  $10^{-5}$



Divergence:  $\sim$  mrad



Output power:



Laser pulse	intensity	Average Power in XUV
10W (our laser)	$10^{-5}$	100 $\mu$ W
2kW (fiber laser)	$10^{-5}$	20mW



# Quasi-Phase matching



Producing XUV depends on phase

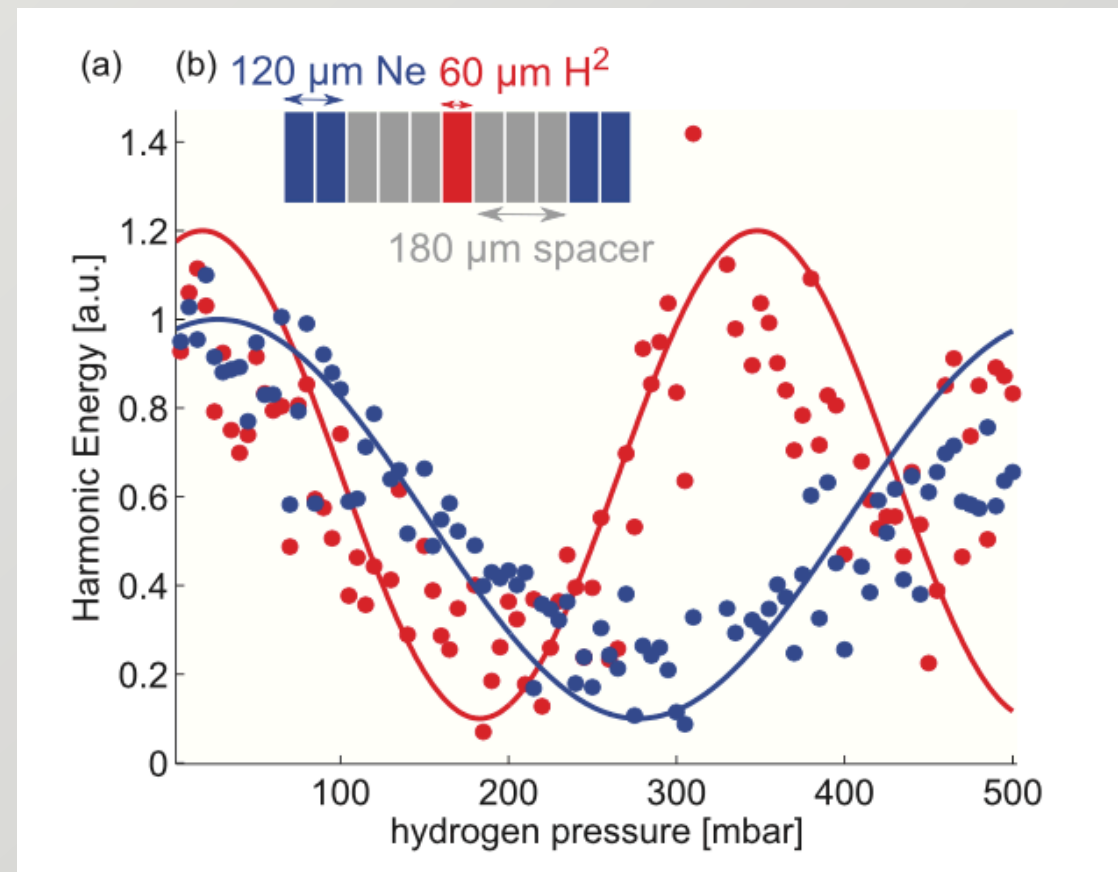
$$\Delta\varphi = \varphi_{\text{fund}} - \varphi_{\text{XUV}} < \pi$$



Dephasing limit for HHG  $\propto l \cdot p$



Reabsorbing XUV



A. Hage,..., M. Wünsche, RSI 103105 (2014)



Overcome: stop reabsorbing by using a phase shifter



## 2. Spectral broadening



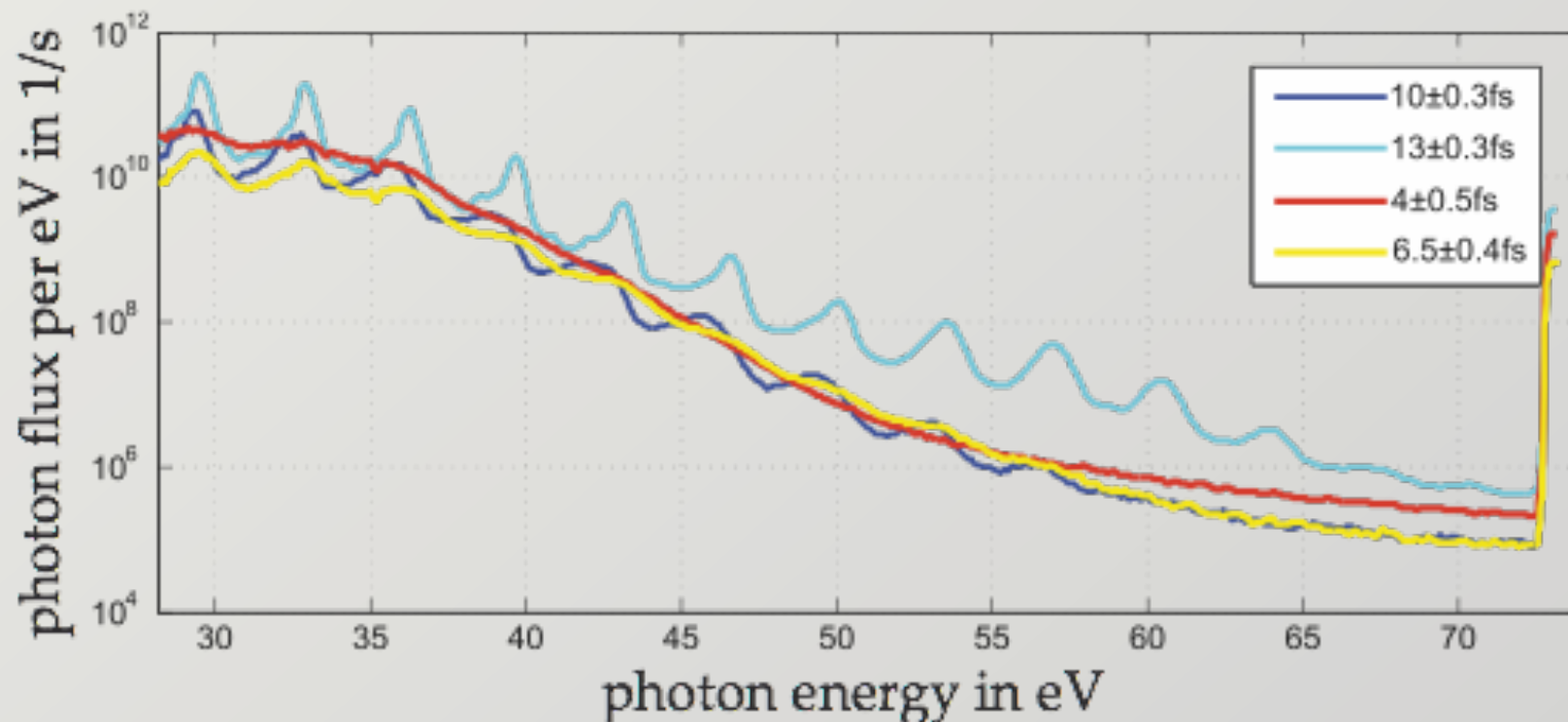
Using few-cycle laser pulse (cycles  $< 3$ )



Reduced number of temporal emitters



Flattening spectral distribution

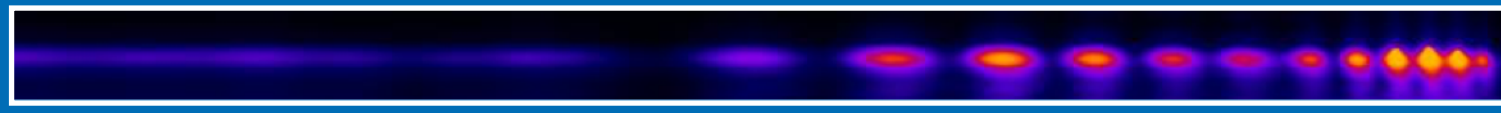


spectrum      photon yield

25...55eV

$10^7 \dots 10^{10}$   
 $\text{eV}^{-1} \text{s}^{-1}$





## Conclusion



HHG: laser-like, coherence, spectral broadness,  
small divergence



Not usable for Lithography



Useful as imaging source



# Field of Application: XUV Coherence tomography

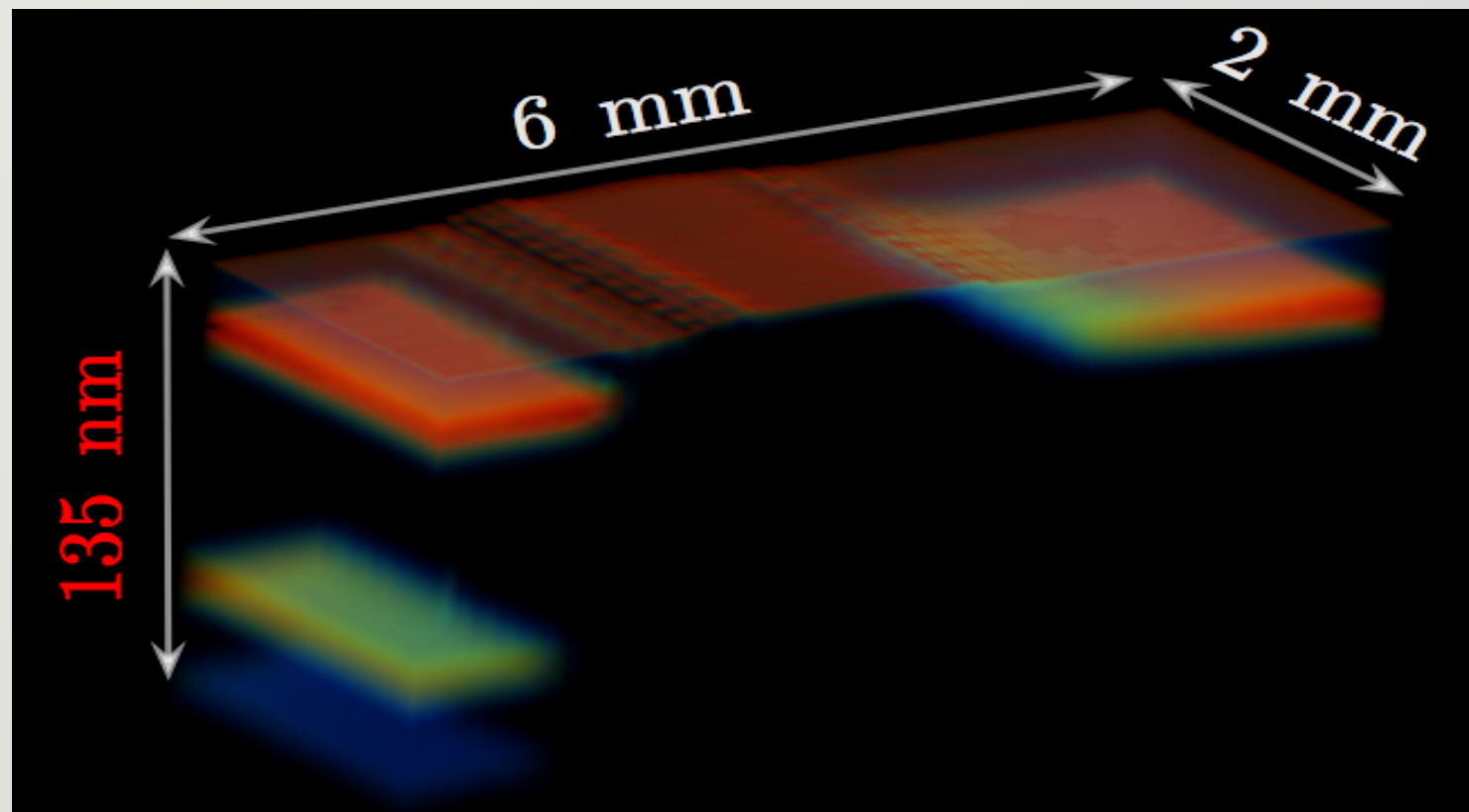


Using broad spectrum of XUV for resolution

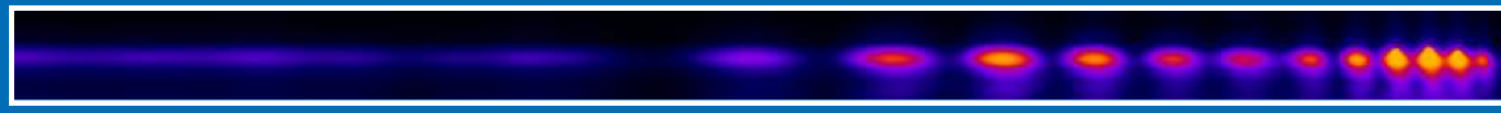


Resolution: few nm

S. Fuchs, ..., M. Wünsche, Appl. Phys. B (2012) 106:789-795



Visit the poster: Nanometer optical coherence tomography using broadband extreme ultra violet light (S44)



**Thank you for  
your attention!**